

Lecture – 19

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- Filter Realization
- Richard's Transformation
- Kuroda Identities
- Stepped Impedance Method



#### **Filter Implementations**

**Q:** So, we now know how to make any and all filters with **lumped** elements but this is a **RF/microwave** engineering course!

You said that lumped elements where difficult to make and implement at high frequencies. You said that distributed elements were used to make microwave components. So how do we make a filter with distributed elements!?!

A: There are **many** ways to make RF/microwave filters with distributed elements. Perhaps the most straightforward is to "**realize**" each individual lumped element with transmission line sections, and then insert these **approximations** in our lumped element solutions.

The **first** of these realizations is: Richard's Transformations

To easily **implement** Richard's Transforms in a microstrip or stripline circuit, we must apply one of **Kuroda's Identities**.



#### **Richard's Transformations**

• Recall the input impedances of short-circuited and open-circuited transmission line **stubs**.



Note that the input impedances are purely **reactive**—just like **lumped** elements!

 However, the reactance of lumped inductors and capacitors have a much different mathematical form to that of transmission line stubs:





 $Z_{in}^s \neq Z_L$ 

 $Z_{in}^o \neq Z_o$ 

# **Richard's Transformations (contd.)**

 In other words, the impedance of transmission line stubs and lumped elements (capacitors and inductors) are different functions with respect to **frequency**. Therefore, we can say in general that, for example:

However, for a given lumped element (L or C) and a given stub (with a given Z<sub>0</sub> and length *l*) the functions **will** be equal at precisely **one frequency**!

For example, there is one frequency—let's call it ω<sub>c</sub> —that satisfies this equation for a given L, Z<sub>0</sub>, and *l*:

$$j\omega_{c}L = jZ_{0}\tan\beta_{c}l = jZ_{0}\tan\left[\frac{\omega_{c}}{v_{p}}l\right]$$

Similarly:  $\begin{bmatrix} -j \\ -j \end{bmatrix}$ 

larly: 
$$\frac{-j}{\omega_c C} = -jZ_0 \cot \beta_c l = -jZ_0 \cot \left\lfloor \frac{\omega_c}{v_p} l \right\rfloor$$

• To make things easier, let's set the **length** of our transmission line stub to  $\lambda_c/8$ , where:

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# **Richard's Transformations (contd.)**

**Q:** Why  $l = \lambda_c / 8$  ?

A: Well, for **one** reason,  $\beta_c l = \pi/4$  and therefore tan  $(\pi/4) = 1.0!$ 

• This greatly **simplifies** our earlier results:

$$j\omega_c L = jZ_0 \tan\left(\frac{\pi}{4}\right) = jZ_0 \qquad \qquad \frac{-j}{\omega_c C} = -jZ_0 \cot\left(\frac{\pi}{4}\right) = -jZ_0$$

• Therefore, if we wish to build a **short-circuited** stub with the **same** impedance as an **inductor** L at frequency  $\omega_c$ , we set the **characteristic impedance** of the stub transmission line to be  $Z_0 = \omega_c L$ :

$$Z_{L} = j\omega_{c}L = Z_{in}^{s}$$

$$Z_{0}, \beta$$

$$Z_{0} = \omega_{c}L$$

$$l = \frac{\lambda_{c}}{8}$$



#### **Richard's Transformations (contd.)**

• Similarly, if we wish to build **open-circuited** stub with the **same** impedance as a **capacitor** C at  $\omega_c$ , we set the **characteristic impedance** of the stub transmission line to be  $Z_0 = \frac{1}{\omega_c C}$ :



We call these two results as **Richard's Transformation**.

However, it is important to remember that Richard's Transformations do **not** result in **perfect** replacements for lumped elements—the stubs **do not** behave like capacitors and inductors!

# **Richard's Transformations (contd.)**

- Instead, the transformation is **perfect**—the impedances are **equal**—at **only one frequency** ( $\omega_c$ ).
- We can use Richard's transformations to replace the inductors and capacitors of a lumped element filter design. In fact, for **low-pass filter design**, the frequency  $\omega_c$  is the filter's **cut-off frequency**.
- Using these stubs to **replace** inductors and capacitors will result in a filter response **similar** to that of the lumped element design—a low pass filter with cut-off frequency  $\omega_c$ .
- However, the behavior of the filter in the **stop-band** will be **very** different from the lumped element design. For example, at the (high) frequencies where the stub length becomes a **multiple** of  $\lambda/2$ , the filter response will be that of  $\omega = 0$ —near perfect **transmission**!

# **Richard's Transformations (contd.)**



**Q:** So **why** does the filter response match the lumped element response so **well** in the **pass-band**?

A: To see why, we first note that the **Taylor Series approximation** for  $tan\varphi$  and  $cot\varphi$  when  $\varphi$  is small (i.e.,  $\varphi \ll 1$ ) is:

$$\begin{array}{c} tan\varphi \approx \varphi \\ \hline and \\ \hline cot\varphi \approx \frac{1}{\varphi} \\ \hline \varphi \ll 1 \end{array}$$

arphi is expressed in **radians**.

 The impedance of Richard's transformation shorted stub at some arbitrary frequency ω is therefore:

$$Z_{in}^{s}(\omega) = jZ_{0} \tan\left(\beta \frac{\lambda_{c}}{8}\right) = j(\omega_{c}L) \tan\left(\frac{\omega}{\omega_{c}} \frac{\pi}{4}\right)$$



## **Richard's Transformations (contd.)**

• Therefore, when  $\omega \ll \omega_c$  (i.e., frequencies in the **pass-band** of a low-pass filter!), we can **approximate** this impedance as:

$$Z_{in}^{s}(\omega) = j(\omega_{c}L)\tan\left(\frac{\omega}{\omega_{c}}\frac{\pi}{4}\right) \approx j(\omega_{c}L)\left(\frac{\omega}{\omega_{c}}\frac{\pi}{4}\right) = j\omega L\left(\frac{\pi}{4}\right)$$
  
Compare this to a  
$$Z_{L} = j\omega L$$
  
Umped inductor  
impedance

Since the value  $\pi/4$  is **relatively** close to one, we find that the Richard's Transformation shorted stub has an input impedance **very** close to the lumped element inductor for **all** frequencies **less than**  $\omega_c$  (i.e., all frequencies of the low-pass filter pass-band)!

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#### **Richard's Transformations (contd.)**

• Similarly, we find that the Richard's transformation **open-circuit** stub, when  $\omega \ll \omega_c$ , has an input impedance of **approximately**:

$$Z_{in}^{o}(\omega) = \frac{-j}{\omega_c C} \cot\left(\frac{\omega}{\omega_c} \frac{\pi}{4}\right) \approx \frac{-j}{\omega_c C} \left(\frac{\omega_c}{\omega} \frac{4}{\pi}\right) = \frac{1}{j\omega C} \left(\frac{4}{\pi}\right)$$
Compare this to a
$$Z_C = \frac{1}{j\omega C}$$
Compare this to a
$$Iumped \text{ capacitor}$$
impedance
e find that results are approximately the **same** for all pass-band

frequencies (i.e., when  $\omega \ll \omega_c$ ).



#### **Kuroda's Identities**

- We will find that **Kuroda's Identities** can be very useful in making the implementation of Richard's transformations more **practicable**.
- Kuroda's Identities essentially provide a list of **equivalent** two port networks. By equivalent, we mean that they have **precisely** the same scattering/impedance/admittance/transmission matrices.
- In other words, we can replace one two-port network with its equivalent in a circuit, and the behavior and characteristics (e.g., its scattering matrix) of the circuit will not change!
- **Q**: Why would we want to do this?
- A: Because one of the equivalent may be more **practical** to implement!

#### For example, we can use Kuroda's Identities to:

- 1. Physically **separate** transmission line stubs.
- 2. Transform series stubs into **shunt** stubs.
- 3. Change impractical **characteristic impedances** into more realizable ones.



- Four Kuroda's identities are provided in a very **ambiguous** and confusing table (Table 8.7) in your **book**. We will find the **first two** identities to be the most useful.
- Consider the following two-port network, constructed with a length of transmission line, and an open-circuit shunt stub:



- The first Kuroda identity states that this two-port network is precisely the same two-port network
- $n^{2} = 1 + \frac{Z_{02}}{Z_{01}} \qquad l \qquad \frac{Z_{01}}{n^{2}}$

as the following:

- Thus,wecanreplacethefirststructureinsomecircuit with this, andthethe behavior of thatcircuitwillnotchangeinthe
- Note this equivalent circuit uses a shortcircuited series stub.



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#### **Kuroda's Identities (contd.)**

 The second of Kuroda's Identities states that this two port network: Is **precisely identical** to this two-port network:



- With regard to Richard's Transformation, these identities are useful when we replace the series inductors with shorted stubs.
- To see why this is useful when implementing a lowpass filter with distributed elements, consider this third order filter example, realized using Richard's Transformations:

Note that we have a few **problems** in terms of implementing this design!





- First of all the stubs are ideally infinitely close to each other how do we build that? We could physically separate them, but this would introduce some transmission line length between them that would mess up our filter response!
- Secondly, series stubs are difficult to construct in microstrip/ stripline—we like shunt stubs much better!
- To solve these problems, we first **add** a short length of transmission line ( $Z_0$  and  $l = \lambda_c/8$ ) to the **beginning** and **end** of the filter:



- Note adding these lengths only results in a phase shift in the filter response—the transmission and reflection functions will remain unchanged.
- Then we can use the second of Kuroda's Identities to replace the series stubs with shunts:



<u>Where:</u>  $n_1^2 = 1 + \frac{Z_0}{\omega_c L_1}$   $n_3^2 = 1 + \frac{Z_0}{\omega_c L_3}$ 

Now **this** is a realizable filter! Note the **three stubs** are separated, and they are all **shunt** stubs.



#### **Stepped-Impedance Low-Pass Filters**

- Another distributed element realization of a lumped element low-pass filter designs is the stepped-impedance low-pass filter.
- These filters are also known as "hi-Z, low-Z" filters, and we're about to find out why!

All distributed elements (e.g., transmission lines, coupled lines, resonators, stubs) exhibit **some** frequency dependency. If we are clever, we can construct these structures in a way that their frequency dependency (i.e.,  $S_{21}(\omega)$ ) conforms to a desirable function of  $\omega$ .

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 $Z_{11} - Z_{21}$ 

#### **Stepped-Impedance Low-Pass Filters (contd.)**

- Say we know the impedance matrix of a **symmetric** two-port device:
- Regardless of the construction of this two port device, we can model it as a simple "T-circuit", consisting of three impedances:
- In other words, if the two series impedances have an impedance value equal to  $Z_{11} Z_{21}$ , and the shunt impedance has a value equal to  $Z_{21}$ , the impedance matrix of this "T-circuit" is:

$$\mathbf{Z} = \begin{bmatrix} Z_{11} & Z_{21} \\ Z_{21} & Z_{11} \end{bmatrix}$$

Thus, **any** symmetric two-port network can be modeled by this "T-circuit"!

$$Z = \begin{bmatrix} Z_{11} & Z_{21} \\ Z_{21} & Z_{11} \end{bmatrix}$$

 $Z_{11} - Z_{21}$ 

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#### **Stepped-Impedance Low-Pass Filters (contd.)**

- For example, consider a length l of transmission line (a symmetric two-port network!):
- (or determine for Recall yourself!) that the impedance parameters of this two port network are:
- With a little **trigonometry**:
- For small  $\beta l$ :

$$Z_{11} - Z_{12} \approx j Z_0 \left(\frac{\beta l}{2}\right)$$

$$Z_{11} - Z_{12} = jZ_0 \tan\left(\frac{\beta l}{2}\right)$$

$$Z_{12} = Z_{21} = -jZ_0 \operatorname{cosec} \beta l \approx \frac{Z_0}{j\beta l}$$



 $Z_0$ 

$$Z_{11} = Z_{22} = -jZ_0 \cot \beta l$$

$$Z_{12} = Z_{21} = -jZ_0 \operatorname{cosec} \beta l$$

$$\sim l \longrightarrow$$

#### **Stepped-Impedance Low-Pass Filters (contd.)**

• Thus, an electrically short  $(\beta l \ll 1)$  transmission line can be **approximately modeled** with a "T-circuit" as:



- Now, consider also the case where the **characteristic impedance** of the transmission line is **relatively large**. We'll **denote** this large characteristic impedance as  $Z_0^h$ .
- Note the **shunt** impedance value  $\frac{Z_0^h}{j\beta l}$ . Since the **numerator**  $(Z_0^h)$  is relatively **large**, and the **denominator**  $(j\beta l)$  is **small**, the impedance of shunt device is **very large**.
- So large, in fact, that we can approximate it as an **open circuit**!

$$\left(\frac{Z_0^h}{j\beta l}\approx\infty\right)$$

For 
$$\beta l \ll 1$$
 and  $Z_0^h \gg Z_0$ 

## Stepped-Impedance Low-Pass Filters (contd.)

 So now we have a further simplification of our model:



 The remaining impedances are now in series, so the circuit can be further simplified to:



• Now, consider the case where the **characteristic impedance** of the transmission line has a relatively **low value**, denoted as  $Z_0^l$ , where  $Z_0^l \ll Z_0$ .

#### Stepped-Impedance Low-Pass Filters (contd.)

• In such a case: 
$$Z_{11} - Z_{12} \approx j Z_0^l \left(\frac{\beta l}{2}\right) \approx 0$$

For  $\beta l \ll 1$  and  $Z_0^l \ll Z_0$ 

• So now we have **another simplification** of our model:



Q: But, what does all this have to do with constructing a low-pass filter?
A: Plenty! Recall that a lossless low-pass filter constructed with lumped elements consists of a "circuit ladder" of series inductors and shunt capacitors!



A: Look at the two equivalent circuits for an electrically short transmission line. The one with large characteristic impedance  $Z_0^h$  has the form of a series inductor, and the one with small characteristic impedance  $Z_0^l$  has the form of a shunt capacitor!



**Q**: Yikes! Our inductance appears to be a function of **frequency**  $\omega$ . I **assume** we simply set this value to cut-off frequency  $\omega_c$ , just like we did for Richard's transformation?

A: Nope! We can simplify the result a bit more. Recall that  $\beta = \omega / v_n$ , so that:

 In other words, the series impedance resulting from our short transmission line is:

**Q: Wow!** This realization seems to give us a result that **precisely** matches an inductor at **all** frequencies—**right**?

A: Not quite! Recall this result was obtained from applying a few approximations—the result is not exact!

## **Stepped-Impedance Low-Pass Filters (contd.)**

• Thus, the "series inductance" of  $L = \frac{Z_0^h \beta l}{\omega}$ 

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$$L = \frac{Z_0^h \beta l}{\omega} = \frac{Z_0^h l}{v_p}$$

$$Z = j\omega L = j\omega \left(\frac{Z_0^h l}{v_p}\right)$$







Moreover, one of the approximations was that the electrical length of the transmission line be small. This obviously cannot be true at all frequencies. As the signal frequency increases, so does the electrical length—our approximate solution will no longer be valid.

Thus, this realization is accurate **only** for "**low** frequencies"— recall that was **likewise** true for **Richard's transformations**!

#### **Q:** Low-frequencies? How **low** is **low**?

A: Well, for our filter to provide a response that **accurately** follows the **lumped element** design, our approximation should be valid for frequencies **up** to (and including!) the filter cut-off frequency  $\omega_c$ .

• A general "rule-of-thumb" is that a small electrical length is defined as being less than  $\pi/4$  radians. Thus, to maintain this small electrical length at frequency  $\omega_c$ , our realization must satisfy the relationship:



- Note that this criterion is difficult to satisfy if the filter cut-off frequency and/or the inductance value L that we are trying to realize is large.
- Our only recourse for these challenging conditions is to increase the value of characteristic impedance  $Z_0^h$ .
- **Q:** Is there some particular difficulty with increasing  $Z_0^h$ ?

A: Could be! There is always a practical limit to how large (or small) we can make the characteristic impedance of a transmission line.

For example, a **large** characteristic impedance implemented in **microstip/stripline** requires a **very narrow** conductor **width W.** But manufacturing tolerances, power handling capability and/or line loss (line resistance R increases as W decreases) place a **lower bound** on how **narrow** we can make these conductors!

However, assuming that we can satisfy the above constraint, we can approximately "realize" a lumped inductor of inductance value L by selecting the correct characteristic impedance Z<sub>0</sub><sup>h</sup> and line length l of our short transmission



**Q:** For **Richard's Transformation**, we **first** set the stub length to a **fixed** value (i.e.,  $l = \lambda_c/8$ ), and **then** determined the **specific characteristic impedance** necessary to realize a **specific inductor value** L. I **assume** we follow the same procedure here?

A: Nope! When constructing stepped-impedance low-pass filters, we typically do the opposite!

**1.** First, we select the value of  $Z_0^h$ , making sure that the short electrical length inequality is **satisfied** for the **largest inductance value** L in our lumped element filter:



 $l_n =$ 

This characteristic impedance value is typically used to realize **all** inductor values L in our low-pass filter, **regardless** of the actual value of inductance L.

2. Then, we determine the **specific lengths**  $l_n$  of the transmission line required to realize **specific** filter **inductors values**  $L_n$ :



- **Q:** What about the **shunt capacitors**?
- A: Almost forgot!
  - Recall the low-impedance transmission line provided a shunt impedance that matched a shunt capacitor:



• Thus, the "**shunt capacitance**" of our transmission line length is:

## **Stepped-Impedance Low-Pass Filters (contd.)**

 And thus the shunt reactance of our transmission line realization is:

$$Z = \frac{-j}{\omega} \left( \frac{v_p Z_0^l}{l} \right)$$

Although this again **appears** to provide **exactly** the same behavior as a **capacitor** (as a function of frequency), it is likewise accurate **only** for **low frequencies**, where  $\beta l < \frac{\pi}{4}$ .

- Thus from our realization **equality**:  $\frac{\beta l}{Z_0^l} = \omega C$
- We can conclude that for our approximations to be valid at all frequencies up to the filter cut-off frequency, the following inequality must be valid:

$$\beta_c l = \omega_c C Z_0^l < \frac{\pi}{4}$$

Note that for **difficult** design cases where  $\omega_c$  and/or C is **very large**, the line **characteristic impedance**  $Z_0^{\ l}$  must be made **very small**.

**Q:** I suppose there is **likewise** a problem with making  $Z_0^l$  very small?

A: Yes! In microstrip and stripline, making  $Z_0^{\ l}$  small means making conductor width W very large. In other words, it will take up lots of space on our substrate. For most applications the surface area of the substrate is both limited and precious, and thus there is generally a practical limit on how wide we can make width W (i.e., how low we can make  $Z_0^{\ l}$ ).



However, assuming that we can satisfy the above constraint, we can approximately "realize" a lumped capacitor of inductance value C by selecting the correct characteristic impedance Z<sub>0</sub><sup>l</sup> and line length l of our short transmission line:



- The **design rules** for **shunt capacitor realization** are:
  - **1.** First, we select the value of  $Z_0^l$ , making sure that the short electrical length inequality is satisfied for the largest capacitance value C in our lumped element filter:



This characteristic impedance value is typically used to realize **all** capacitor values C in our low-pass filter, **regardless** of the actual value of capacitance C.

2. Then, we determine the **specific lengths**  $l_n$  of the transmission line required to realize **specific** filter capacitor values  $C_n$ :

$$l_n = \left(v_p Z_0^l\right) C_n$$



 An example of a low-pass, stepped-impedance filter design is provided on page 414-416 of your book

